

## An Approach to the Schrödinger Equation

### ***Classical Mechanics: Energy Conservation:***

$$E = \frac{mv^2}{2} + V = \frac{p^2}{2m} + V \quad \text{with} \quad p = mv$$

### ***Quantum physics: particles have wavelength and angular frequency:***

$$p = \frac{h}{\lambda} \quad E_{\text{kin}} = \hbar\omega \quad \hbar = \frac{h}{2\pi} \quad \omega = 2\pi f$$

### ***A deBroglie particle-wave:***

$$\Psi = A \exp \left[ i \left( \frac{2\pi}{\lambda} x - \omega t \right) \right] = A \exp \left[ \frac{i}{\hbar} (px - E_{\text{kin}} t) \right]$$

Here  $E_{\text{kin}}$  is the kinetic energy of the moving particle!

### ***Quantum Mechanics: The Schrödinger Equation***

Differentiating the particle wave equation twice with respect to x leads to

$$\frac{d^2}{dx^2} \Psi = -\frac{p^2}{\hbar^2} \Psi$$

Differentiating it once with respect to t and replacing  $E_{\text{kin}}$  with E from classical mechanics leads to

$$\frac{d}{dt} \Psi = -\frac{i}{\hbar} E \Psi = -\frac{i}{\hbar} \left( \frac{p^2}{2m} + V \right) \Psi$$

This is a generalization since the  $E_{\text{kin}}$  in the deBroglie wave is actually just the kinetic energy. We are looking for an equation that is also true for particles in a potential so we replaced the kinetic energy with the total energy. Generalizations like this lead to a new theory which cannot be derived from any other theory. It is not possible to derive a theory, one can only develop it. One theory never follows from another. If it did it wouldn't be a theory.

From the two last equations we can construct a differential equation as follows:

We multiply the first by  $-\hbar^2/2m$  and the second by  $i\hbar$  to give:

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \Psi = \frac{p^2}{2m} \Psi$$

And

$$i\hbar \frac{d}{dt} \Psi = \left( \frac{p^2}{2m} + V \right) \Psi$$

If we add  $V\Psi$  to the first of these equations the right hand sides are identical so the left hand sides should be too and we get a differential equation from the left hand sides.

$$i\hbar \frac{d}{dt} \Psi = \left[ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V \right] \Psi$$

This is the **time-dependent Schroedinger equation in one dimension**. It can be generalized to three dimensions with the abbreviation

$$\nabla^2 = \frac{d^2}{dx^2} + \frac{d^2}{dy^2} + \frac{d^2}{dz^2}$$

To give the **time-dependent Schroedinger equation**.

$$i\hbar \frac{d}{dt} \Psi = \left[ -\frac{\hbar^2}{2m} \nabla^2 + V \right] \Psi$$

The expression in brackets is called the **Hamilton operator  $H$**  (it performs a mathematical operation on  $\Psi$ ).

$$i\hbar \frac{d}{dt} \Psi = H\Psi$$

$\Psi$  is called the **wave function**. It is a complex valued function of  $x, y, z$  and  $t$ . The wave function has no physical interpretation but its **absolute square  $|\Psi|^2$**  is understood as the **probability** to find the particle inside a small volume at position  $x, y, z$  and time  $t$ , divided by the size of this volume. This is a **probability density**.

If a particle is bound in a potential, for example an electron bound in an atom, the probability density does not depend on time, The probability to find a particle in a certain volume in a certain position is always the same. The orbital does not change shape.

A wave function with a time-independent probability density can be constructed as follows:

$$\Psi(x, y, z, t) = \exp\left(-i \frac{E}{\hbar} t\right) \psi(x, y, z)$$

with an absolute square that depends only on  $x, y$  and  $z$  and not on time since

$$|\Psi|^2 = \Psi \Psi^* = \exp\left(-i \frac{E}{\hbar} t\right) \psi(x, y, z) \exp\left(i \frac{E}{\hbar} t\right) \psi^*(x, y, z) = \psi \psi^* = |\psi|^2$$

We'll insert this wave function into the time-dependent Schrödinger equation now but we will stick to the one-dimensional case

$$\Psi(x, t) = \exp\left(-i \frac{E}{\hbar} t\right) \psi(x)$$

to give:

$$i\hbar \frac{d}{dt} \left[ \exp\left(-i \frac{E}{\hbar} t\right) \psi(x) \right] = \left[ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V \right] \left[ \exp\left(-i \frac{E}{\hbar} t\right) \psi(x) \right]$$

The differentiation with respect to t on the left can be carried out directly, whereas on the right the exponential can be moved in front of the bracket since it only depends on t and not on x.

$$i\hbar \left(-i \frac{E}{\hbar}\right) \exp\left(-i \frac{E}{\hbar} t\right) \psi(x) = \exp\left(-i \frac{E}{\hbar} t\right) \left[ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V \right] \psi(x)$$

which simplifies to

$$E \psi(x) = \left[ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V \right] \psi(x)$$

or swapping right and left side, using  $H$  for the Hamilton operator and generalizing to three dimensions

$$H \psi(x, y, z) = E \psi(x, y, z)$$

This is the **time-independent Schrödinger equation** which describes orbitals that are constant in time.