

First we define known quantities

$$\begin{aligned}
 \Delta &:= r^2 - 2 M r + a^2 \\
 \Sigma &:= r^2 + a^2 \cos^2[\theta] \\
 J &:= C \text{DiracDelta}[r - r_0] \text{DiracDelta}\left[\theta - \frac{\pi}{2}\right] \\
 J_{m\uparrow} &:= \left( \sqrt{2} \left( r - i a \cos[\theta] \right) \right)^{-1} + i (r^2 + a^2) \sin[\theta] J \\
 J_n &:= - \frac{a \Delta}{\Sigma} \sin[\theta]^2 J
 \end{aligned}$$

Now for  $J_2$ , we can ignore  $\partial_\phi$  since  $J_m$  and  $J_n$  are independent of it

$$\begin{aligned}
 J_2 = & \frac{-\Delta}{2\sqrt{2} \Sigma (r - i a \cos \theta)^2} \left[ \sqrt{2} \left( \frac{\partial}{\partial r} - \frac{a}{\Delta} \frac{\partial}{\partial \phi} + \frac{1}{r - i a \cos \theta} \right) (r - i a \cos \theta)^2 J_{\bar{m}} \right. \\
 & \left. + 2 \left( \frac{\partial}{\partial \theta} - \frac{i}{\sin \theta} \frac{\partial}{\partial \phi} + \frac{i a \sin \theta}{r - i a \cos \theta} \right) \frac{\Sigma (r - i a \cos \theta)}{\Delta} J_n \right] \quad (2.12)
 \end{aligned}$$

$$\begin{aligned}
 J_2 := & \frac{-\Delta}{2 \sqrt{2} \Sigma (r - i a \cos[\theta])^2} \\
 & \left( \sqrt{2} \left( \partial_r \# + \frac{1}{r - i a \cos[\theta]} \# \right) \&@ \left( (r - i a \cos[\theta])^2 J_{m\uparrow} \right) + 2 \left( \partial_\theta \# + \frac{i a \sin[\theta]}{r - i a \cos[\theta]} \# \right) \&@ \right. \\
 & \left. \left( \frac{\Sigma (r - i a \cos[\theta])}{\Delta} J_n \right) \right)
 \end{aligned}$$

Finally for the integration we need to first look at the definition of the SpinWeighted spherical harmonics.

Since all SpinWeighted spherical harmonics, here denoted  ${}_s Y_{lm}(\theta, \phi) = Y[s, l, \theta, \phi]$  behave in the  $\phi$  argument as  $\text{Exp}[i m \phi]$  and since all other variables are independent of  $\phi$  in the equation above We have

$\int Y[s, l, m, \theta, \phi] = 2 \pi \text{DiracDelta}[m, 0] Y[s, l, m, \theta, 0]$ . Then we can integrate over  $\theta$ .

the source term is given by

$${}^2 J_{lm}(r) = \int_0^{2\pi} \int_0^\pi \frac{(r - i a \cos \theta)^2}{(r_+ - r_-)^2} \Sigma J_2 {}_{-1} \bar{Y}_{lm} \sin \theta d\theta d\phi$$

In[ ]:=

We can split the integrand into two parts depending if we have  $\text{DiracDelta}\left[\theta - \frac{\pi}{2}\right]$  or the derivative  $\text{DiracDelta}'\left[\theta - \frac{\pi}{2}\right]$ .

and we will use the standard properties

$$\int_{-\infty}^{\infty} f(x) \delta(x-a) dx = f(a)$$

$$\int f(x) \delta^{(n)}(x) dx \equiv - \int \frac{\partial f}{\partial x} \delta^{(n-1)}(x) dx.$$

First part

$$\begin{aligned} \text{In}[*]:= & \frac{(r - i a \cos[\theta])^2}{(rp - rm)^2} \Sigma J2 \text{Conjugate}[Y[-1, l, 0, \theta, 0]]; \\ & \% /. \text{DiracDelta}[-\pi + 2 \theta] \rightarrow 1 /. \text{DiracDelta}[-\pi + 2 \theta] \rightarrow 0; \\ & \text{FirstPart} = \% /. \theta \rightarrow \frac{\pi}{2} \\ \text{Out}[*]= & -\frac{1}{2 \sqrt{2} (-rm + rp)^2} (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]] \\ & \left( -8 i a^2 C \text{DiracDelta}[r - r0] + \sqrt{2} \left( 3 r \left( \frac{1}{\sqrt{2} r} + 2 i C (a^2 + r^2) \text{DiracDelta}[r - r0] \right) + \right. \right. \\ & \left. \left. r^2 \left( -\frac{1}{\sqrt{2} r^2} + 4 i C r \text{DiracDelta}[r - r0] + 2 i C (a^2 + r^2) \text{DiracDelta}'[r - r0] \right) \right) \right) \end{aligned}$$

Second part

$$\begin{aligned} \text{In}[*]:= & \frac{(r - i a \cos[\theta])^2}{(rp - rm)^2} \Sigma J2 \text{Conjugate}[Y[-1, l, 0, \theta, 0]]; \\ & \% /. \text{DiracDelta}[-\pi + 2 \theta] \rightarrow 0 /. \text{DiracDelta}[-\pi + 2 \theta] \rightarrow 1; \\ & \text{SecondPart} = -D[\%, \theta] /. \theta \rightarrow \frac{\pi}{2} /. \text{Conjugate}'[Y[-1, l, 0, \frac{\pi}{2}, 0]] \rightarrow 1 \\ \text{Out}[*]= & -\frac{2 i \sqrt{2} a^2 C (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]] \text{DiracDelta}[r - r0]}{(-rm + rp)^2} + \\ & \frac{(a^2 - 2 M r + r^2) (2 - 8 a C r \text{DiracDelta}[r - r0]) Y^{(0,0,0,1,0)}[-1, l, 0, \frac{\pi}{2}, 0]}{2 \sqrt{2} (-rm + rp)^2} \end{aligned}$$

We get the solution

In[ ]:= FirstPart + SecondPart // Simplify

$$\begin{aligned} \text{Out[ ]} = & - \frac{2 i \sqrt{2} a^2 C (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]] \text{DiracDelta}[r - r_0]}{(-r m + r p)^2} - \\ & \frac{1}{2 \sqrt{2} (-r m + r p)^2} (a^2 - 2 M r + r^2) \text{Conjugate}[Y[-1, l, 0, \frac{\pi}{2}, 0]] \\ & \left( -8 i a^2 C \text{DiracDelta}[r - r_0] + \sqrt{2} \left( 3 r \left( \frac{1}{\sqrt{2} r} + 2 i C (a^2 + r^2) \text{DiracDelta}[r - r_0] \right) + \right. \right. \\ & \quad \left. \left. r^2 \left( -\frac{1}{\sqrt{2} r^2} + 4 i C r \text{DiracDelta}[r - r_0] + 2 i C (a^2 + r^2) \text{DiracDelta}'[r - r_0] \right) \right) \right) + \\ & \frac{(a^2 - 2 M r + r^2) (2 - 8 a C r \text{DiracDelta}[r - r_0]) Y^{(0,0,0,1,0)}[-1, l, 0, \frac{\pi}{2}, 0]}{2 \sqrt{2} (-r m + r p)^2} \end{aligned}$$

We should further utilize identities for the Y function but we can already see that we must have made a mistake

Since the solution given in the article ( $C = l \left( \frac{\Delta_0}{r_0} \right)$ ) and ( $e = 0$ )

$$\begin{aligned} {}^2J_{lm} = & - \frac{\Delta \delta_{m0}}{\sqrt{2} (r_+ - r_-)^2} [(M a e / \mathfrak{Q}_0) + \pi \mathfrak{J}(\Delta_0 / \mathfrak{Q}_0)^{1/2}] \\ & \cdot \left[ i (r_0^2 + a^2) {}_{-1}\bar{Y}_{l0} \left( \frac{\pi}{2}, 0 \right) \delta'(r - r_0) \right. \\ & \left. + \left\{ i r_0 {}_{-1}\bar{Y}_{l0} \left( \frac{\pi}{2}, 0 \right) - a [l(l+1)]^{1/2} {}_0\bar{Y}_{l0} \left( \frac{\pi}{2}, 0 \right) \right\} \delta(r - r_0) \right] \end{aligned}$$

Is only quadratic in r, but our solution is at least quartic in r.